

on the other hand, if $[\hat{A}, \hat{B}] \neq 0$
 results from \hat{C} depends on whether or not
 \hat{B} measurement has actually been performed.

3) Uncertainty:

* D14 operator $\Delta\hat{A} \equiv \hat{A} - \langle\hat{A}\rangle\mathbb{1}$ $g(x) = \frac{1}{\sigma\sqrt{2\pi}} e^{-\frac{1}{2}\frac{(x-\mu)^2}{\sigma^2}}$

$$\begin{aligned} \langle(\Delta\hat{A})^2\rangle &: \text{dispersion of } \hat{A} \quad (\sim \sigma^2) \\ &= \langle(\hat{A} - \langle\hat{A}\rangle\mathbb{1})^2\rangle \\ &= \langle\hat{A}^2 - 2\hat{A}\langle\hat{A}\rangle + \langle\hat{A}\rangle^2\mathbb{1}\rangle \\ &= \langle\hat{A}^2\rangle - \underline{2\langle\hat{A}\rangle\langle\hat{A}\rangle} + \underline{\langle\hat{A}\rangle^2} \\ &= \langle\hat{A}^2\rangle - \langle\hat{A}\rangle^2 \end{aligned} \quad \left. \begin{array}{l} \text{variance} \\ \text{or mean squared} \\ \text{deviation.} \end{array} \right\}$$

(std deviation $\sigma = \sqrt{\langle\hat{A}^2\rangle - \langle\hat{A}\rangle^2}$)

$\langle(\Delta\hat{A})^2\rangle$ vanish iff when state $|\psi\rangle$ in question.

is an eigenstate of \hat{A}

Proof: $\hat{A}|\psi\rangle = a|\psi\rangle$

$$\begin{aligned} \rightarrow \Delta\hat{A} \left(\hat{A} - \langle\hat{A}\rangle\mathbb{1} \right) |\psi\rangle &= a|\psi\rangle - \underbrace{\langle\psi|\hat{A}|\psi\rangle}_{a|\psi\rangle} \mathbb{1} |\psi\rangle \\ &= a|\psi\rangle - a|\psi\rangle = 0 \end{aligned}$$

uncertainty principle:

normalized.

Given two Herm. operator. \hat{A} & \hat{B} , and a state $|\psi\rangle$

$$\langle (\Delta \hat{A})^2 \rangle \langle (\Delta \hat{B})^2 \rangle \geq \frac{1}{4} |\langle [\hat{A}, \hat{B}] \rangle|^2$$

To proof. we need Lemmas (Intermediate theorems)

Lemma 1: Schwartz inequality:

$$\langle \alpha | \alpha \rangle \langle \beta | \beta \rangle \geq |\langle \alpha | \beta \rangle|^2$$

$$|\vec{a}|^2 |\vec{b}|^2 \geq |\vec{a} \cdot \vec{b}|^2$$

↓
 $|\vec{a}| |\vec{b}| \cos \theta$

Lemma 2: The expectation value of \hat{X}_H is real.

Lemma 3: The expectation value of Anti-Herm.
is purely imaginary.

$$\hat{X}^\dagger = -\hat{X} \quad \hat{X}|\psi\rangle = a|\psi\rangle$$

$$\rightarrow \langle \psi | \frac{\hat{X}|\psi\rangle}{a|\psi\rangle} = a \quad \xrightarrow{a = -a^*} \quad a^* \langle \psi | \psi \rangle$$

$$= \langle \psi | (-\hat{X}^\dagger) | \psi \rangle = -\langle \hat{X} \psi | \psi \rangle = -\langle a \psi | \psi \rangle$$

w/ Lemma 1 Let $|\alpha\rangle = \Delta \hat{A} | \rangle$
 \hookrightarrow may be any ket

$$|\beta\rangle = \Delta \hat{B} | \rangle$$

Based on $\langle \alpha | \alpha \rangle \langle \beta | \beta \rangle \geq |\langle \alpha | \beta \rangle|^2$

$$\begin{array}{ccc}
 \langle \Delta \hat{A} | \Delta \hat{A} \rangle & \langle \Delta \hat{B} | \Delta \hat{B} \rangle & |\langle \Delta \hat{A} | \Delta \hat{B} \rangle|^2 \\
 \downarrow & \downarrow & \downarrow \\
 \langle \Delta \hat{A}^2 \rangle & \langle \Delta \hat{B}^2 \rangle & \langle \Delta \hat{A} \Delta \hat{B} \rangle
 \end{array}$$

$$\Delta \hat{A} \Delta \hat{B} = \frac{1}{2} [\Delta \hat{A}, \Delta \hat{B}] + \frac{1}{2} \{ \Delta \hat{A}, \Delta \hat{B} \}$$

$$\begin{array}{cc}
 \frac{\Delta \hat{A} \Delta \hat{B} - \Delta \hat{B} \Delta \hat{A}}{2} & \frac{\Delta \hat{A} \Delta \hat{B} + \Delta \hat{B} \Delta \hat{A}}{2}
 \end{array}$$

$$\begin{aligned}
 ([\Delta \hat{A}, \Delta \hat{B}])^\dagger &= (\Delta \hat{A} \Delta \hat{B} - \Delta \hat{B} \Delta \hat{A})^\dagger \\
 &= (\Delta \hat{B} \Delta \hat{A} - \Delta \hat{A} \Delta \hat{B}) \\
 &= -([\Delta \hat{A}, \Delta \hat{B}])
 \end{aligned}$$

Previously given.

$$(\hat{X} \hat{Y})^\dagger = \hat{Y}^\dagger \hat{X}^\dagger$$

$\Rightarrow [\Delta \hat{A}, \Delta \hat{B}]$ is anti-Herm.

likewise $\{ \Delta \hat{A}, \Delta \hat{B} \}$ is a Herm.

$$\rightarrow \langle \Delta \hat{A} \Delta \hat{B} \rangle = \frac{1}{2} \langle [\Delta \hat{A}, \Delta \hat{B}] \rangle + \frac{1}{2} \langle \{ \Delta \hat{A}, \Delta \hat{B} \} \rangle$$

\downarrow imaginary
 \downarrow real

$$\rightarrow |\langle \Delta \hat{A} \Delta \hat{B} \rangle|^2 = \frac{1}{4} |\langle [\Delta \hat{A}, \Delta \hat{B}] \rangle|^2 + \frac{1}{4} |\langle \{ \Delta \hat{A}, \Delta \hat{B} \} \rangle|^2$$

since $[\Delta \hat{A}, \Delta \hat{B}] = [\hat{A}, \hat{B}]$ trivial proof. ≥ 0

\downarrow
 $\hat{A} - \langle \hat{A} \rangle \mathbb{I}$

$$\langle \Delta \hat{A} \Delta \hat{A} \rangle \langle \Delta \hat{B} \Delta \hat{B} \rangle \geq |\langle \Delta \hat{A} \Delta \hat{B} \rangle|^2$$



$$\langle (\Delta \hat{A})^2 \rangle \langle (\Delta \hat{B})^2 \rangle \geq \frac{1}{4} |\langle [\Delta \hat{A}, \Delta \hat{B}] \rangle|^2 + \frac{1}{4} |\langle \{\Delta \hat{A}, \Delta \hat{B}\} \rangle|^2$$

$$\geq \frac{1}{4} |\langle [\hat{A}, \hat{B}] \rangle|^2$$

$$\geq 0$$

$$\Rightarrow \langle (\Delta \hat{A})^2 \rangle \langle (\Delta \hat{B})^2 \rangle \geq \frac{1}{4} |\langle [\hat{A}, \hat{B}] \rangle|^2 \quad *$$

e.g. $\hat{A} = \hat{x}$, $\hat{B} = \hat{p}$ $\rightarrow [\hat{x}, \hat{p}] = i\hbar \mathbb{1}$

$$\Rightarrow \underbrace{\langle (\Delta \hat{x})^2 \rangle}_{\delta_x^2} \underbrace{\langle (\Delta \hat{p})^2 \rangle}_{\delta_p^2} \geq \frac{1}{4} \underbrace{|\langle i\hbar \mathbb{1} \rangle|^2}_{\frac{1}{4} \hbar^2}$$

$$\Rightarrow \delta_x^2 \cdot \delta_p^2 \geq \frac{1}{4} \hbar^2 \Rightarrow \delta_x \delta_p \geq \frac{1}{2} \hbar$$

3. Position, momentum & translation.

3-1 Continuous spectra


\hat{x} , & \hat{p} are associated w/ continuous eigenvalue (spectra) e.g. \hat{p}_z eigenvalue can

$$\hat{W} |w\rangle = w |w\rangle$$

continuous spectra

be anywhere $[-\infty, \infty]$

$$g(x) = \frac{1}{\sigma\sqrt{\pi}} e^{-\frac{x^2}{\sigma^2}}$$

$$\sigma \rightarrow 0 \rightarrow \delta(x)$$


$$\hat{A} |a\rangle = a |a\rangle \rightarrow \text{discrete spectra}$$

Analogy:

Discrete.

$$\langle a' | a'' \rangle = \delta_{a'a''}$$

$$\sum_{a'} |a'\rangle \langle a'| = \mathbb{1}$$

$$|\alpha\rangle = \sum_{a'} |a'\rangle \langle a' | \alpha \rangle$$

$$\sum_{a'} |\langle a' | \alpha \rangle|^2 = 1$$

$$\langle a'' | \hat{A} | a' \rangle = a' \delta_{a'a''}$$

$$a' |a'\rangle \begin{cases} \text{if } a' \neq a'' \rightarrow 0 \\ \text{if } a' = a'' \rightarrow \neq 0 \end{cases}$$

$\delta(x)$: Dirac delta function

$\rightarrow 0$ anywhere, except "at" $x=0$

continuous spectra.

$$\langle w' | w'' \rangle = \delta(w' - w'')$$

$$\int dw' |w'\rangle \langle w'| = \mathbb{1}$$

$$|\alpha\rangle = \int dw' |w'\rangle \langle w' | \alpha \rangle$$

$$\int dw' |\langle w' | \alpha \rangle|^2 = 1$$

$$\langle w'' | \hat{W} | w' \rangle = w' \delta(w' - w'')$$

$$\begin{cases} \text{if } w' \neq w'' \rightarrow 0 \\ \text{if } w' = w'' \rightarrow \neq 0 \end{cases}$$