

⇒ i.e. the amplitude for a particle to go from (x', t') to (x'', t'') is sum over all possible intermediate positions x'' at t''

6-3 Path Integral

consider transition from (x_1, t_1) to (x_N, t_N)

$$t_j - t_{j-1} = \Delta t = \frac{t_N - t_1}{N-1}$$

$$\Rightarrow \langle x_N, t_N | x_1, t_1 \rangle$$

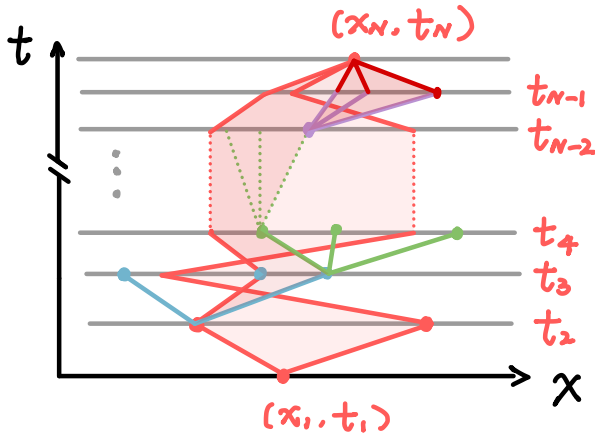
$$K(x_N, t_N; x_{N-1}, t_{N-1})$$

$$= \int dx_{N-1} \int dx_{N-2} \cdots \int dx_2 \langle x_N, t_N | x_{N-1}, t_{N-1} \rangle$$

$$\langle x_{N-1}, t_{N-1} | x_{N-2}, t_{N-2} \rangle$$

⋮

$$\langle x_2, t_2 | x_1, t_1 \rangle$$



$$\Rightarrow K(x_N, t_N; x_1, t_1) = \lim_{N \rightarrow \infty} \int dx_2 \int dx_3 \cdots \int dx_{N-1} \prod_{j=2}^N K(x_j, t_j; x_{j-1}, t_{j-1})$$

path integral is constructed by composing infinitesimally small segment of propagators over time steps.

6-4 A more precise formulation

(Feynman Path Integral Formulation for QM)

⇒ Classic Mechanics based on Lagrangian function.

$$\mathcal{L}(x, \dot{x}, t) = \frac{1}{2} m \dot{x}^2 - V(x)$$

↑
 $x = x(t), \dot{x} = \dot{x}(t)$

Hamilton's principle (principle of stationary action)

⇒ Given a fixed starting and ending space-time point (x_1, t_1) & (x_N, t_N) , Nature will choose the path, so that the action is stationary (typically, the action is minimized)

Now, def. - action: action S is defined as

$$S(N, 1) = \int_{t_1}^{t_N} dt \mathcal{L}(x, \dot{x}, t)$$



⇒ S is sort of the total "effort" nature has put

⇒ if S is stationary ⇒ $\delta S = 0$

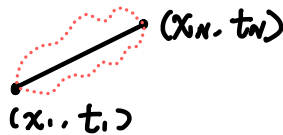
$$\Rightarrow \delta \int_{t_i}^{t_f} dt \mathcal{L}(x, \dot{x}, t) = 0$$

typically, S is minimized.

\Rightarrow Nature is "efficient", not wasting energy.

\Rightarrow NOTE: Newton's Law is just a special case.
& can be derived from this (Hamilton's principle)

e.g., for a free particle, i.e., $V(x) = 0$



*
 \Rightarrow classical mechanics is contrasting Q.M.
in the way that the path is definite in space-time plane, while in Q.M. All possible paths play roles.

\Rightarrow Question: How to reconcile classical mechanics w/ Q.M. in a smooth manner, in the limit of $\hbar \rightarrow 0$
(recall Ehrenfest's Theorem, likewise, this is a reasonable question.)

\Rightarrow Feynman's Question based on Dirac's NOTE.

*

$$e^{i \int_{t_1}^{t_2} \frac{dt L(x, \dot{x}, t)}{\hbar}$$

corresponds to $\langle x_2, t_2 | x_1, t_1 \rangle$?

↓
seed for Feynman's path integral and much of modern Quantum Theory.

what did P. Dirac mean above? classical mechanics.

Now consider: $S(n, n-1) = \int_{t_{n-1}}^{t_n} dt L(x, \dot{x}, t)$

↓
small segment (prescribed)

⇒ then we're instructed by Dirac to associate this segment with $e^{\frac{i}{\hbar} S(n, n-1)}$

$$\begin{aligned} \text{Now, from } 1 \rightarrow N &\Leftrightarrow \prod_{n=2}^N e^{\frac{i}{\hbar} S(n, n-1)} \\ &= e^{\frac{i}{\hbar} \underbrace{\sum_{n=2}^N S(n, n-1)}_{\text{total action from } 1 \rightarrow N}} \\ &= e^{\frac{i}{\hbar} S(N, 1)} \end{aligned}$$

NOTE. This is NOT $\langle x_N, t_N | x_1, t_1 \rangle$ yet.

* This is just one particular path contributing to it.

⇒ then, the total contribution:

$$\langle x_N, t_N | x_1, t_1 \rangle \underset{? \text{ All paths}}{\sim} \sum e^{\frac{i}{\hbar} S(N,1)}$$

Now, if $\hbar \rightarrow 0$

⇒ small change in $S \rightarrow$ big change in $e^{\frac{i}{\hbar} S}$

⇒ $e^{\frac{i}{\hbar} S}$ will oscillate violently.

⇒ most of paths do not contribute anymore
Due to interference.

⇒ Except for the one : $\delta S(N,1) = 0$

