

⇒ General form of unitary operator:

$$\hat{U}(\lambda) = e^{-i\lambda \hat{G}}$$

General.

⇒ it involves an arbitrary operator \hat{G}

⇒ it preserves the probability of a state.

$$\hat{U}^\dagger \hat{U} = \mathbb{1}$$

$$e^{-i\lambda \hat{G}} = 1 + (-i\lambda \hat{G}) + \dots$$

knowing $\hat{T}(dx) = \mathbb{1} - i\hat{K}dx$

$$\Rightarrow \hat{T}(dx) = e^{-idx \hat{K}} = e^{-idx \hat{K}'/\hbar} \quad (\hbar \hat{K}' = \hat{K})$$

$$\begin{aligned} \Rightarrow \hat{T}(dx) &= 1 - i \frac{dx}{\hbar} \hat{K}' + \mathcal{O}(dx^2) \\ &\approx 1 - i \frac{dx}{\hbar} \hat{K}' \end{aligned}$$

Comparing: $\hat{T}(dx) \approx 1 - \frac{d}{dx} dx$

$$\Rightarrow 1 - \frac{d}{dx} dx = 1 - i \frac{dx}{\hbar} \hat{K}' \Rightarrow$$

$$\hat{K}' = -i\hbar \frac{d}{dx}$$

$$\Rightarrow \hat{T}(dx) = \mathbb{1} - i \hat{K} dx = \mathbb{1} - i \frac{\hat{K}'}{\hbar} dx$$

$$\Rightarrow \hat{T}(dx) = \mathbb{1} - i \frac{\hat{P}}{\hbar} dx$$

$$\Rightarrow \hat{P} = \hbar \hat{K}$$

↓
 \hat{P}
↓

Dimension analysis.

$$\hbar: [J \cdot s] = [J \cdot Hz^{-1}]$$

Energy per Hz.

$$\hbar/dx : \underline{[J \cdot s \cdot cm^{-1}]}$$

↳ unit for momentum.

$$\text{unit for } \hat{k} \text{ is } \frac{[J \cdot s \cdot cm^{-1}]}{[J \cdot s]} = cm^{-1}$$

$$\Rightarrow \hat{T}(dx') = \mathbb{1} - \frac{i}{\hbar} \hat{p} dx'$$

$$\begin{aligned} \Rightarrow \underline{[\hat{x}_i, \hat{p}_j]} &= \hat{x}_i \hat{p}_j - \hat{p}_j \hat{x}_i = \hbar \hat{x}_i \hat{k}_j - \hbar \hat{k}_j \hat{x}_i \\ &= \hbar [\hat{x}_i, \hat{k}_j] = \underline{\mathbb{1} \cdot i \cdot \hbar \cdot \delta_{ij}} \end{aligned}$$

⇒ Position momentum uncertainty:

$$\langle (\Delta X)^2 \rangle \langle (\Delta P)^2 \rangle \geq \frac{\hbar^2}{4}$$

Now translate over $\Delta X'$

$$dx' = \lim_{N \rightarrow \infty} \frac{\Delta X'}{N}$$

$$\rightarrow \hat{T}(\Delta X' \cdot \underset{\substack{\downarrow \\ \text{unit vector.}}}{\vec{x}}) = \lim_{N \rightarrow \infty} \left(\mathbb{1} - \frac{i}{\hbar} \hat{p} \cdot \frac{\Delta X'}{N} \right)^N$$

unit vector.

$$e = \lim_{n \rightarrow \infty} \left(1 + \frac{x}{n} \right)^n$$

$$\Rightarrow e^x = \lim_{n \rightarrow \infty} \left(1 + \frac{x}{n} \right)^n$$

$$\ln \left(1 + \frac{x}{n} \right)^n = n \ln \left(1 + \frac{x}{n} \right)$$

$$\lim_{n \rightarrow \infty} n \ln \left(1 + \frac{x}{n} \right) \approx n \cdot \frac{x}{n} = x$$

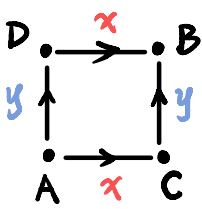
$$\Rightarrow e^{\lim_{n \rightarrow \infty} \ln \left(1 + \frac{x}{n} \right)^n} = e^x \Rightarrow e^x = \lim_{n \rightarrow \infty} \left(1 + \frac{x}{n} \right)^n$$

$$= \lim_{N \rightarrow \infty} \left(1 + \frac{x}{N} \right)^N$$

$$= e^x = e^{-\frac{i}{\hbar} \hat{p} \cdot \Delta X'}$$

"x"

$$= \lim_{N \rightarrow \infty} \left(1 + \frac{-(i/\hbar) \hat{p} \cdot \Delta X'}{N} \right)^N$$



$$\Rightarrow \hat{T}(\Delta y \cdot \vec{y}) \cdot \hat{T}(\Delta x \cdot \vec{x})$$

$$= \hat{T}(\Delta x \cdot \vec{x}) \cdot \hat{T}(\Delta y \cdot \vec{y})$$

$$\Rightarrow [\hat{T}(\Delta y \cdot \vec{y}), \hat{T}(\Delta x \cdot \vec{x})] = 0$$

$$= \left[\left(\mathbb{1} - \frac{i \hat{P}_y \Delta y}{\hbar} - \frac{\hat{P}_y^2 (\Delta y)^2}{2\hbar^2} \dots \right), \left(\mathbb{1} - \frac{i \hat{P}_x \Delta x}{\hbar} - \frac{\hat{P}_x^2 (\Delta x)^2}{2\hbar^2} \dots \right) \right]$$

$$\Leftarrow - \frac{(\Delta x)(\Delta y) [\hat{P}_y, \hat{P}_x]}{\hbar^2} = 0$$

$$\Rightarrow [\hat{P}_y, \hat{P}_x] = 0$$

$$\Rightarrow \text{Generally, } [\hat{P}_i, \hat{P}_j] = 0^*$$

$\Rightarrow \{ \hat{P}_i \}$ form an Abelian Group: Group describes a system where observables commute.

$$* \text{ since } [\hat{T}, \hat{P}] = 0 \text{ and } \hat{P} |P_i\rangle = P_i |P_i\rangle$$

$$\Rightarrow |P_i\rangle \text{ is eigenket of } \hat{T}$$

3-5. canonical commutation:

Fundamental quantum behavior of conjugate variables

$$[\hat{x}_i, \hat{x}_j] = 0$$

$$[\hat{p}_i, \hat{p}_j] = 0$$

$$[\hat{x}_i, \hat{p}_j] = i\hbar \delta_{ij} \mathbb{1}$$

4. wavefunctions in position, & momentum space.

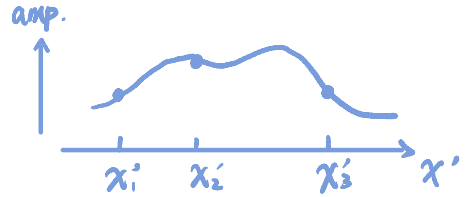
4-1. wavefunctions in position space.

$$\hat{x}|\alpha\rangle = x'|\alpha\rangle \quad \neq 0 \text{ iff } x'' = x'$$

orthogonality: $\langle x''|\alpha\rangle = \delta(x'' - x')$

$$|\alpha\rangle = \int dx' |\alpha\rangle \langle x'|\alpha\rangle$$

\downarrow
 $f_\alpha(x')$



a smooth function, $\psi_\alpha(x')$, in the 1D space.

(the points in the 1D. are labeled by x'), in this

space, $\psi_\alpha(x')$ is a smooth function of x' .

and is called a wavefunction in x' space.

$$\Rightarrow \underline{\psi_\alpha(x') = \langle x'|\alpha\rangle}$$

consistent w/ position $\left\{ \begin{array}{l} \rightarrow \text{Probability amplitude for} \\ \text{finding a particle in state } |\alpha\rangle \end{array} \right.$

$$P = |\langle x'|\alpha\rangle|^2 = |\psi_\alpha(x')|^2 = \psi_\alpha^*(x')\psi_\alpha(x')$$